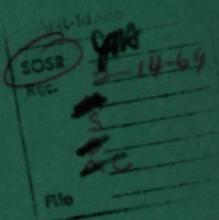


Argonne National Laboratory

**ENDF/B Neutron Cross-section Data
for Natural Helium
(ENDF-125)**

by

E. M. Pennington



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E. M. Pennington

Reactor Physics Division

October 1968

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ENDF/B Neutron Cross-section Data
for Natural Helium
(ENDF-125)

by

E. M. Pennington

ABSTRACT

This report describes the compilation of neutron cross-section data for natural helium. The elastic scattering cross section, $\bar{\mu}_L$, ξ , γ , and elastic scattering Legendre coefficients are included for ${}^4\text{He}$. Since ${}^3\text{He}$ has a natural abundance of merely 0.00013%, only its (n,p) cross section, which is large at low energies, is considered. A listing of the data in the ENDF/B format is presented, as are graphs of much of the data.

I. INTRODUCTION

Since natural helium is a possible coolant for fast breeder reactors, the present helium cross-section data were compiled. The data will be distributed with ENDF/B identification number MAT = 1088.

The composition of natural helium is 99.99987% ${}^4\text{He}$ and 0.00013% ${}^3\text{He}$. Because of the low abundance of ${}^3\text{He}$, only its (n,p) cross section, which is large at low energies, need be considered. Elastic scattering is the only possible reaction for neutrons incident on ${}^4\text{He}$ at energies below 15 MeV. Thus the elastic scattering cross section and values of $\bar{\mu}_L$, ξ , and γ are included in the compilation, as well as elastic scattering Legendre coefficients. Here $\bar{\mu}_L$ is the average cosine of the angle of elastic scattering in the laboratory system, ξ is the average logarithmic energy decrement in an elastic collision, and γ , the Greuling-Goertzel parameter, is the average of the square of the logarithmic energy decrement divided by twice the average logarithmic energy decrement. Parameters for a free-gas thermal scattering law are also included.

II. DATA SOURCES AND CALCULATIONAL METHODS

The elastic scattering cross section and Legendre expansion coefficients were calculated from s-, p-, and d-wave phase shifts using a FORTRAN program written for the purpose. Appendix A presents the equations involved in this program. The phase shifts were read from smooth

curves based on Table I of Ref. 1. At energies below the 300-keV lower limit of this table, each of the two p-wave phase shifts was obtained by using a functional form based on the low-energy limit for a single p-wave resonance, with parameters determined from fitting the low-energy phase shifts of the table. Appendix B presents the equations connected with the low-energy phase shifts, and the parameters obtained by the fitting procedure. The s-wave phase shift below 300 keV was calculated using hard-sphere scattering and a nuclear radius of $a = 2.4$ F. This yields a thermal scattering cross section of $4\pi a^2 = 0.7238$ barn, in agreement with the experimental value² of 0.73 ± 0.05 barn. The low-energy s-wave phase shifts of Ref. 1 are consistent with a nuclear radius of about 2.48 F and so would yield a somewhat high thermal cross section.

Table I lists the phase shifts used in the calculations at energies above 10 keV.

TABLE I. Helium Phase Shifts Used in Calculations of Elastic Scattering Cross Sections and Legendre Expansion Coefficients

E_{lab} , MeV	δ_0^+	δ_1^+	δ_1^-	δ_2^+	δ_2^-	E_{lab} , MeV	δ_0^+	δ_1^+	δ_1^-	δ_2^+	δ_2^-
0.02	3.08203	0.00092	0.00022			2.9	2.4155	2.1660	0.4782	0.0030	0.0021
0.03	3.06865	0.00171	0.00040			3.0	2.4051	2.1677	0.5009	0.0031	0.0023
0.05	3.04742	0.00373	0.00086			3.2	2.3806	2.1712	0.5428	0.0037	0.0026
0.07	3.03017	0.00629	0.00143			3.4	2.3579	2.1729	0.5847	0.0042	0.0031
0.10	3.00842	0.0110	0.00246			3.6	2.3370	2.1694	0.6248	0.0049	0.0035
0.15	2.97849	0.0212	0.00457			3.8	2.3178	2.1625	0.6667	0.0054	0.0038
0.20	2.95326	0.0343	0.00711			4.0	2.2986	2.1537	0.7086	0.0061	0.0044
0.25	2.93103	0.0504	0.0100			4.2	2.2794	2.1468	0.7435	0.0070	0.0049
0.30	2.9025	0.0698	0.0133			4.4	2.2602	2.1398	0.7767	0.0077	0.0054
0.35	2.8833	0.0930	0.0170			4.6	2.2427	2.1328	0.8081	0.0084	0.0059
0.40	2.8658	0.1204	0.0209			4.8	2.2253	2.1258	0.8360	0.0093	0.0065
0.45	2.8466	0.1536	0.0252			5.0	2.2078	2.1171	0.8622	0.0101	0.0070
0.50	2.8292	0.1920	0.0299			5.2	2.1904	2.1084	0.8849	0.0110	0.0075
0.55	2.8152	0.2356	0.0348			5.4	2.1747	2.0996	0.9058	0.0120	0.0082
0.60	2.8047	0.2845	0.0401			5.6	2.1590	2.0892	0.9250	0.0131	0.0087
0.65	2.7908	0.3421	0.0454			5.8	2.1450	2.0804	0.9442	0.0141	0.0094
0.70	2.7786	0.4067	0.0524			6.0	2.1293	2.0717	0.9617	0.0152	0.0103
0.75	2.7646	0.4800	0.0576			6.2	2.1153	2.0630	0.9774	0.0162	0.0108
0.80	2.7524	0.5725	0.0646			6.4	2.1014	2.0543	0.9896	0.0175	0.0115
0.85	2.7402	0.6615	0.0698			6.6	2.0857	2.0438	1.0001	0.0185	0.0124
0.90	2.7297	0.7662	0.0768			6.8	2.0717	2.0351	1.0105	0.0197	0.0133
0.95	2.7175	0.8779	0.0838			7.0	2.0577	2.0263	1.0210	0.0209	0.0140
1.00	2.7070	0.9983	0.0908			7.2	2.0455	2.0176	1.0280	0.0222	0.0145
1.05	2.6965	1.1170	0.0977			7.4	2.0316	2.0089	1.0350	0.0236	0.0152
1.10	2.6861	1.2235	0.1065			7.6	2.0193	2.0001	1.0420	0.0248	0.0159
1.13	2.6791	1.2881	0.1100			7.8	2.0054	1.9914	1.0472	0.0262	0.0166
1.15	2.6756	1.3299	0.1134			8.0	1.9932	1.9827	1.0507	0.0279	0.0175
1.17	2.6721	1.3701	0.1169			8.2	1.9809	1.9740	1.0559	0.0297	0.0183
1.20	2.6669	1.4312	0.1222			8.4	1.9705	1.9652	1.0594	0.0314	0.0190
1.25	2.6564	1.5202	0.1292			8.6	1.9583	1.9565	1.0612	0.0332	0.0197
1.30	2.6477	1.5970	0.1379			8.8	1.9478	1.9478	1.0629	0.0349	0.0206
1.35	2.6389	1.6668	0.1466			9.0	1.9356	1.9391	1.0647	0.0367	0.0215
1.40	2.6302	1.7296	0.1553			9.2	1.9234	1.9321	1.0647	0.0384	0.0223
1.45	2.6215	1.7802	0.1641			9.4	1.9129	1.9234	1.0647	0.0401	0.0234
1.5	2.6128	1.8239	0.1745			9.6	1.9007	1.9146	1.0647	0.0419	0.0243
1.6	2.5953	1.9059	0.1937			9.8	1.8902	1.9059	1.0629	0.0436	0.0253
1.7	2.5796	1.9652	0.2129			10.0	1.8780	1.8972	1.0629	0.0454	0.0262
1.8	2.5639	2.0106	0.2321			10.5	1.8535	1.8780	1.0612	0.0506	0.0288
1.9	2.5482	2.0438	0.2531			11.0	1.8291	1.8570	1.0577	0.0541	0.0314
2.0	2.5342	2.0752	0.2758	0.0012	0.0009	11.5	1.8047	1.8378	1.0507	0.0593	0.0340
2.1	2.5203	2.0961	0.2967	0.0012	0.0009	12.0	1.7802	1.8169	1.0402	0.0646	0.0367
2.2	2.5063	2.1153	0.3194	0.0014	0.0010	12.5	1.7593	1.7994	1.0297	0.0698	0.0384
2.3	2.4923	2.1293	0.3403	0.0016	0.0012	13.0	1.7383	1.7820	1.0193	0.0750	0.0419
2.4	2.4784	2.1398	0.3630	0.0017	0.0012	13.5	1.7174	1.7628	1.0071	0.0803	0.0436
2.5	2.4644	2.1468	0.3857	0.0019	0.0014	14.0	1.6965	1.7453	0.9966	0.0885	0.0454
2.6	2.4522	2.1537	0.4084	0.0021	0.0016	14.5	1.6755	1.7279	0.9861	0.0908	0.0471
2.7	2.4400	2.1590	0.4328	0.0024	0.0017	15.0	1.6546	1.7104	0.9756	0.0960	0.0489

Values of $\bar{\mu}_L$, ξ , and γ were calculated from the Legendre coefficients using a FORTRAN program, MUXIGA. This program is based on the equations derived in Refs. 3-5. These equations are summarized in Appendix C.

An elastic scattering Legendre-coefficient transformation matrix from the center-of-mass system to the laboratory system was computed using CHAD.⁶ A minor revision was made in the CHAD code to avoid elements of the matrix being set equal to zero if they were larger than elements closer to the diagonal.

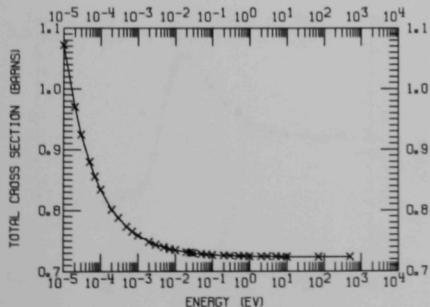
The (n,p) cross section for ${}^3\text{He}$ is that recommended in the ${}^3\text{He}$ evaluation by J. Als-Nielsen.⁷ Extension from 10 to 15 MeV was made using linear extrapolation on a log σ versus log E scale.

The total cross section is the sum of the elastic scattering and (n,p) cross sections.

The free-gas thermal scattering law involves a value of 0.7238 barn for the free-atom scattering cross section.

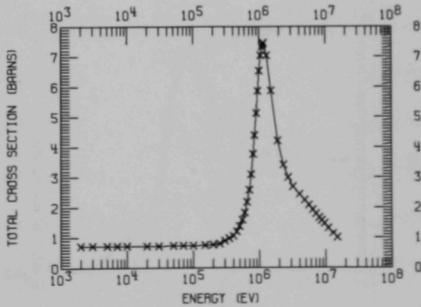
III. GRAPHS AND LISTING

Figures 1-8 are graphs of the ENDF/B File 3 data.⁸ The File 3 data are represented by a series of pairs of x and y values with an interpolation rule applying between successive pairs. The interpolation rule might specify, for example, that $\log y$ is linear in $\log x$ between successive (x,y) pairs. In Figs. 1-8 the axis types have been chosen to correspond to the interpolation rule. Thus in Figs. 1 and 2 a linear axis is used for the total cross section and a logarithmic axis is used for the energy, because linear y versus $\log x$ interpolation is specified for this cross section. The total cross section



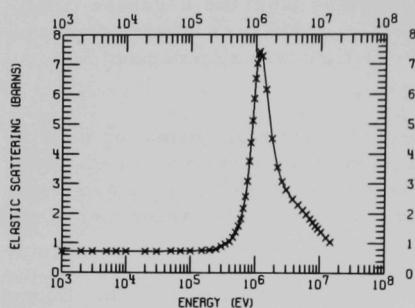
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Fig. 1. Helium Total Cross Section,
 10^{-5} eV to 1 keV



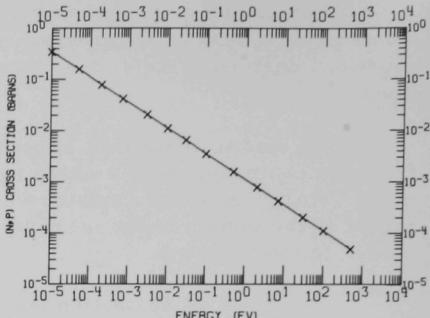
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Fig. 2. Helium Total Cross Section,
1 keV to 15 MeV

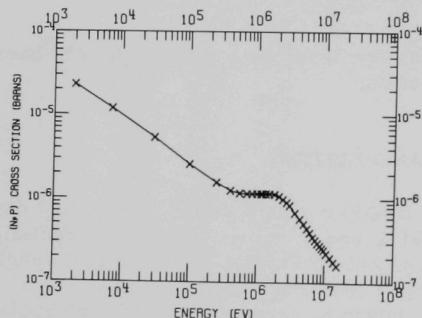


113-681

Fig. 3. Helium Elastic Scattering Cross Section

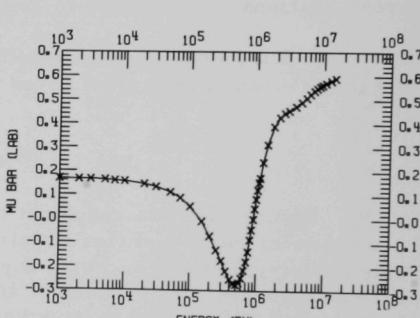


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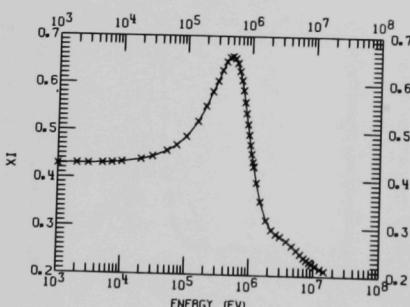
Fig. 4. Helium (n,p) Cross Section, 10^{-5} eV to 1 keV

113-685

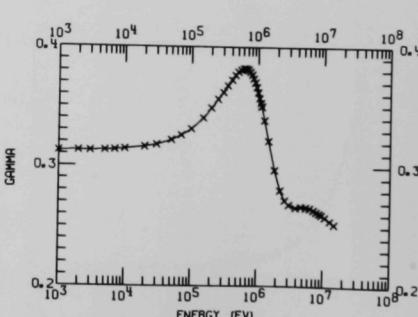
Fig. 5. Helium (n,p) Cross Section, 1 keV to 15 MeV



113-696

Fig. 6. Helium Average Elastic Scattering Cosine, $\bar{\mu}_L$ 

113-698

Fig. 7. Helium Average Logarithmic Energy Decrement, ξ 

113-680

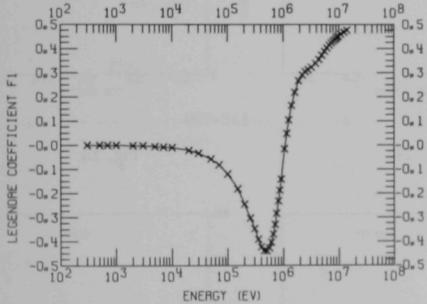
Fig. 8. Helium Greuling-Goertzel Parameter, γ

and the (n,p) cross section have been plotted in two energy ranges in Figs. 1 and 2 and Figs 4 and 5, respectively, in order to display the data better. The data in Figs. 3, 6, 7, and 8 for the elastic scattering cross section, $\bar{\mu}_L$, ξ , and γ are not plotted below 1 keV because the values are almost constant at low energies. Not all data points are plotted, especially at high energies, to avoid clutter on the graphs.

Figures 9-12 display the Legendre coefficients for elastic scattering angular distributions in the center-of-mass system, f_1 through f_4 . The elastic scattering angular distribution, $d\sigma(E)/d\Omega$, is given by⁸

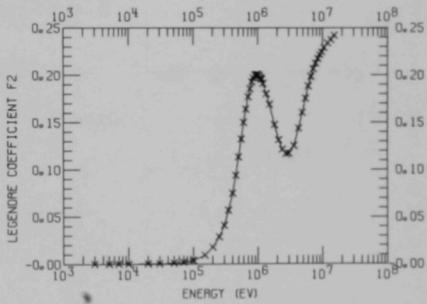
$$\frac{d\sigma(E)}{d\Omega} = \frac{\sigma_s(E)}{4\pi} \sum_{\ell=0}^{\infty} (2\ell+1) f_\ell(E) P_\ell(\mu), \quad (1)$$

where the elastic scattering cross section $\sigma_s(E)$ is tabulated in ENDF/B File 3, and the expansion coefficients $f_\ell(E)$ are presented at a series of



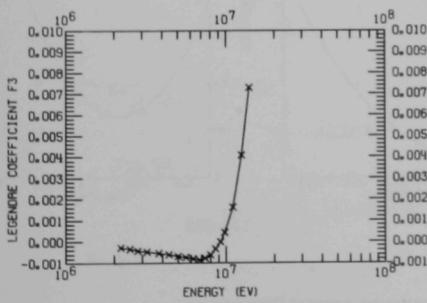
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Fig. 9. Helium Elastic Scattering
Legendre Coefficient f_1



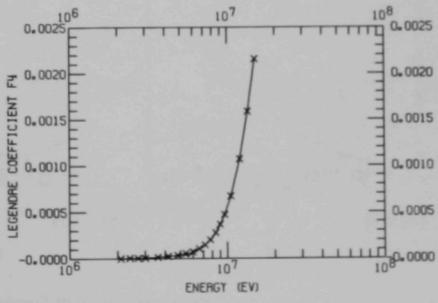
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Fig. 10. Helium Elastic Scattering
Legendre Coefficient f_2



113-702

Fig. 11. Helium Elastic Scattering
Legendre Coefficient f_3

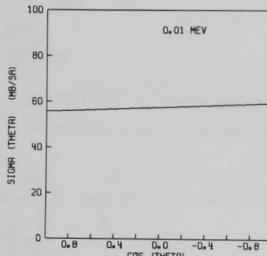


113-674

Fig. 12. Helium Elastic Scattering
Legendre Coefficient f_4

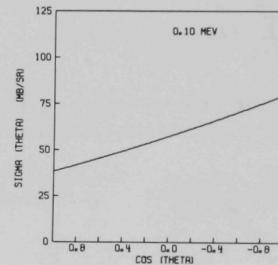
energies in File 4. In Figs. 13-31 the angular distributions have been calculated from Eq. 1 at a representative sequence of 19 energies. Note that the angular distributions at the energies shown below 1.2 MeV are peaked in the backward direction, as is typical on the low-energy side of a p-wave resonance. All graphs were drawn with the aid of the CDC-3600 computer and the Calcomp 580 plotter.

The helium data in the ENDF/B format are listed in Appendix D. The listing will not be explained here since all information necessary for understanding the formats involved is in Ref. 8.



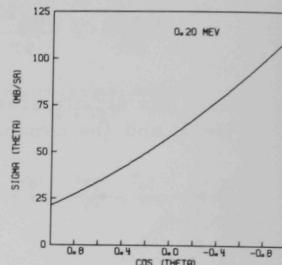
113-701

Fig. 13



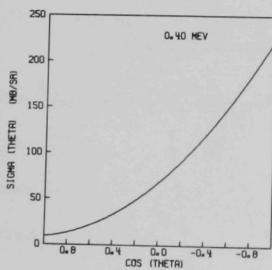
113-678

Fig. 14



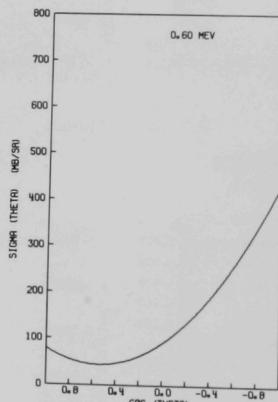
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Fig. 15



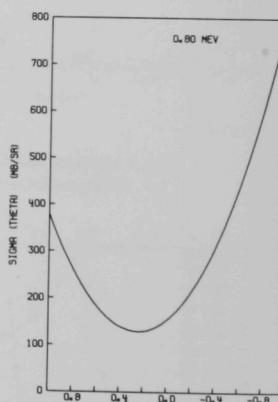
113-676

Fig. 16



113-703

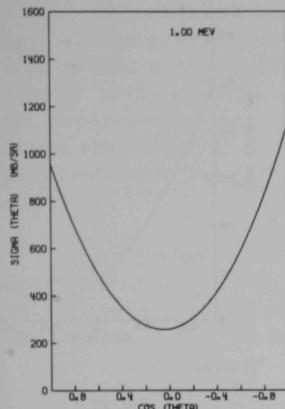
Fig. 17



113-683

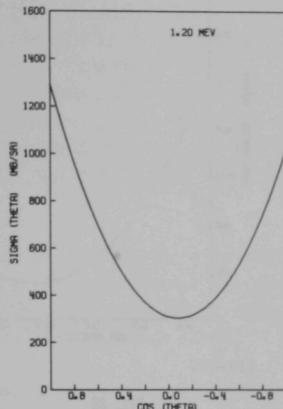
Fig. 18

Figs. 13-18. Helium Elastic Scattering Angular Distributions



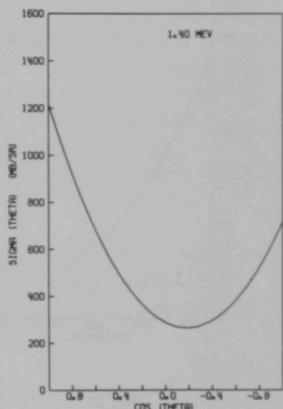
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Fig. 19



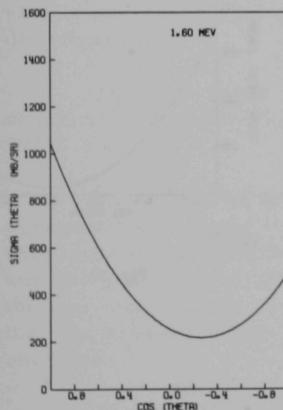
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Fig. 20



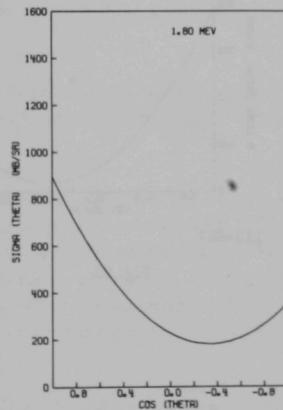
113-699

Fig. 21



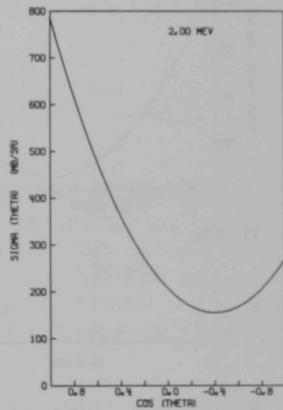
113-695

Fig. 22



113-684

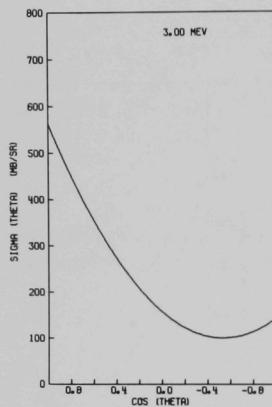
Fig. 23



113-692

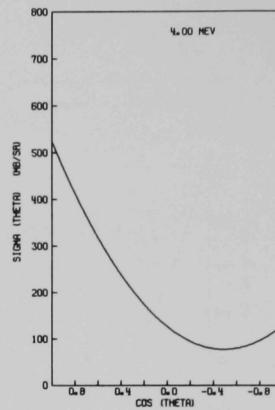
Fig. 24

Figs. 19-24. Helium Elastic Scattering Angular Distributions

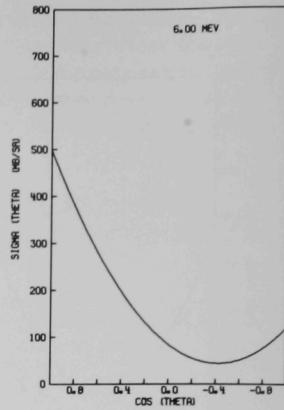


113-693

Fig. 25

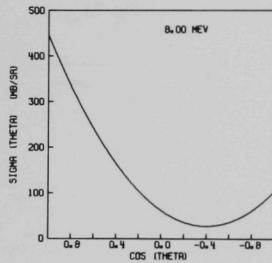


113-675



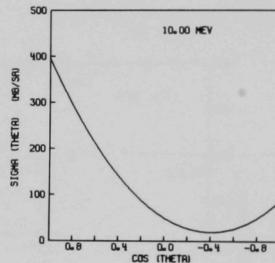
113-690

Fig. 27



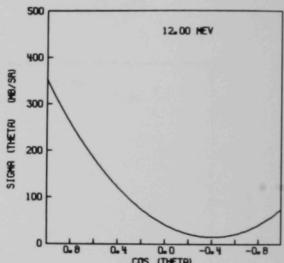
113-679

Fig. 28



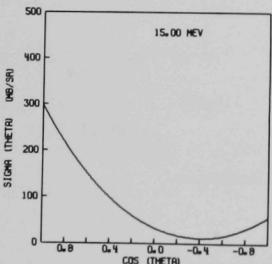
113-687

Fig. 29



113-704

Fig. 30



113-689

Fig. 31

Figs. 25-31

Helium Elastic Scattering Angular Distributions

IV. DISCUSSION

The phase shifts of Ref. 1, which were used in the compilation, are optical-model phase shifts chosen to fit both angular-distribution and polarization data at many energies. The total cross section is also fit within the scatter of the experimental points. Because Ref. 1 is recent and considers the experimental data comprehensively, the data calculated from the phase shifts should be reliable. Another recent set of phase shifts⁹ is not very different from that of Ref. 1 and could also have been used in the present work.

As discussed in Ref. 7, the ^3He (n,p) cross section is rather well known. Probably more error is introduced into the (n,p) cross section for natural helium by the uncertainty in the ^3He isotopic abundance than by the uncertainty in the ^3He (n,p) cross section itself.

Previous evaluations of natural helium for reactor calculations include those of Schmidt¹⁰ and of Buckingham *et al.*^{11,12} Schmidt's evaluation includes the (n,p) cross section for ^3He and σ_s , $\bar{\mu}_L$, and a set of phase shifts for ^4He . Buckingham *et al.* present separate compilations for ^3He and ^4He . For ^3He , elastic, (n,p), (n,d), and ($n,2n$) cross sections are given, as well as elastic angular distributions. The ^4He evaluation consists of σ_s and angular distributions.

The (n,p) cross section of ^3He , which was based on Ref. 7, is not very different from those of previous evaluations,^{10,12} as may be seen from Table 3 of Ref. 7.

The values of the elastic scattering cross section are in fairly good agreement with those in the other evaluations.^{10,11} Schmidt's values¹⁰ agree with those of the present evaluation at energies below the main resonance and are generally a little larger at higher energies, except for a range on the high-energy side of the main resonance. The English values¹¹ are higher than the present values at energies below about 1 MeV and lower at higher energies.

The present values of $\bar{\mu}_L$ are in rather good agreement with those shown on Fig. He-3 in Part I of KFK-120¹⁰ at energies above approximately 1 MeV, but do not agree well at lower energies.

APPENDIX A

Equations Relating the Elastic Scattering Cross Section and Legendre Expansion Coefficients to the s-, p-, and d-wave Phase Shifts

The elastic scattering angular distribution, $d\sigma/d\Omega$, in the center-of-mass system is given by¹³

$$\frac{d\sigma}{d\Omega} = \frac{1}{4k^2} (|a|^2 + |b|^2) = \sum_n C_n P_n(\mu), \quad (\text{A.1})$$

where

$$a = \sum_{\ell=0}^{\infty} \left[(\ell+1)(e^{2i\delta_{\ell}^+} - 1) + \ell(e^{2i\delta_{\ell}^-} - 1) \right] P_{\ell}(\mu), \quad (\text{A.2})$$

and

$$b = \sum_{\ell=0}^{\infty} \left(e^{2i\delta_{\ell}^+} - e^{2i\delta_{\ell}^-} \right) P_{\ell}^1(\mu). \quad (\text{A.3})$$

Here μ is the cosine of the center-of-mass scattering angle, $P_{\ell}(\mu)$ is the Legendre polynomial of order ℓ , and

$$P_{\ell}^1(\mu) = (1 - \mu^2)^{1/2} \frac{d}{d\mu} P_{\ell}(\mu).$$

The quantity $k = \sqrt{2mE/\hbar^2}$, where E is the energy of relative motion of neutron and nucleus in the center-of-mass system, and m is the reduced mass. The phase shifts δ_{ℓ}^+ and δ_{ℓ}^- refer to $j = \ell + 1/2$ and $j = \ell - 1/2$, respectively, where j is the angular momentum quantum number referring to the addition of the orbital and spin angular momenta.

On restricting the sums in Eqs. A.2 and A.3 by considering only s-, p-, and d-wave phase shifts, and using the explicit forms for the Legendre polynomials, we obtain

$$k^2 C_0 = \sin^2 \delta_0^+ + 2 \sin^2 \delta_1^+ + \sin^2 \delta_1^- + 3 \sin^2 \delta_2^+ + 2 \sin^2 \delta_2^-, \quad (\text{A.4})$$

$$k^2 C_1 = 3 \sin^2 \delta_0^+ + 6 \sin^2 \delta_1^+ + 3 \sin^2 \delta_1^- + \frac{18}{5} \sin^2 \delta_2^+ + \frac{12}{5} \sin^2 \delta_2^-$$

$$- 2 \sin^2 (\delta_0^+ - \delta_1^+) - \sin^2 (\delta_0^+ - \delta_1^-) - \frac{18}{5} \sin^2 (\delta_1^+ - \delta_2^+)$$

$$- 2 \sin^2 (\delta_1^- - \delta_2^-) - \frac{2}{5} \sin^2 (\delta_1^+ - \delta_2^-), \quad (\text{A.5})$$

$$\begin{aligned}
 k^2 C_2 = & 5 \sin^2 \delta_0^+ + 4 \sin^2 \delta_1^+ + 2 \sin^2 \delta_1^- + \frac{51}{7} \sin^2 \delta_2^+ + \frac{34}{7} \sin^2 \delta_2^- \\
 & - 2 \sin^2 (\delta_1^+ - \delta_1^-) - \frac{6}{7} \sin^2 (\delta_2^+ - \delta_2^-) - 3 \sin^2 (\delta_0^+ - \delta_2^+) \\
 & - 2 \sin^2 (\delta_0^+ - \delta_2^-),
 \end{aligned} \tag{A.6}$$

$$\begin{aligned}
 k^2 C_3 = & 6 \sin^2 \delta_1^+ + 3 \sin^2 \delta_1^- + \frac{27}{5} \sin^2 \delta_2^+ + \frac{18}{5} \sin^2 \delta_2^- \\
 & - \frac{12}{5} \sin^2 (\delta_1^+ - \delta_2^+) - \frac{18}{5} \sin^2 (\delta_1^+ - \delta_2^-) - 3 \sin^2 (\delta_1^- - \delta_2^+),
 \end{aligned} \tag{A.7}$$

and

$$k^2 C_4 = \frac{54}{7} \sin^2 \delta_2^+ + \frac{36}{7} \sin^2 \delta_2^- - \frac{36}{7} \sin^2 (\delta_2^+ - \delta_2^-). \tag{A.8}$$

Equation 1 gives the angular distribution in the ENDF/B normalization. Thus the elastic scattering cross section is $\sigma_s = 4\pi C_0$, and the Legendre coefficients are given by

$$f_\ell = \frac{C_\ell}{(2\ell+1) C_0}.$$

APPENDIX B

Low-energy Phase Shifts from Resonance Parameters

The phase shift, $\delta_{\ell J}$, near an isolated purely scattering resonance with orbital angular momentum ℓ and total angular momentum J is¹⁴

$$\delta_{\ell J} = \tan^{-1} \left(\frac{R_{\ell J} P_{\ell}}{1 - R_{\ell J} \hat{S}_{\ell}} \right) - \phi_{\ell}, \quad (\text{B.1})$$

where

$$\phi_{\ell} = \tan^{-1} \left(\frac{F_{\ell}}{G_{\ell}} \right) \quad (\text{B.2})$$

with

$$F_{\ell} = \rho j_{\ell}(\rho) \quad (\text{B.3})$$

and

$$G_{\ell} = -\rho n_{\ell}(\rho). \quad (\text{B.4})$$

The functions $j_{\ell}(\rho)$ and $n_{\ell}(\rho)$ are spherical Bessel functions.¹⁵ The quantity ρ is defined by

$$\rho = ka = a \sqrt{2mE/\hbar^2} = a \sqrt{\frac{2m}{\hbar^2} \frac{A}{A+1} E_{\text{lab}}} = C \sqrt{E_{\text{lab}}},$$

where E is the energy of the relative motion of neutron and target in the center-of-mass system, E_{lab} is the neutron energy in the laboratory system, $m = m_n A / (1 + A)$, with A being the ratio of the target mass to the neutron mass, m_n , and a is the nuclear radius. The penetration and shift factors are

$$P_{\ell} = \frac{\rho}{F_{\ell}^2 + G_{\ell}^2} \quad (\text{B.5})$$

and

$$S_{\ell} = \frac{\rho(F_{\ell}' F_{\ell} + G_{\ell}' G_{\ell})}{F_{\ell}^2 + G_{\ell}^2} \quad (\text{B.6})$$

where the prime signifies differentiation with respect to ρ . The \hat{S}_{ℓ} in Eq. B.1 is related to S_{ℓ} by $\hat{S}_{\ell} \equiv S_{\ell} - B_{\ell J}$, where $B_{\ell J}$ is a boundary-condition constant.¹⁴ The function $R_{\ell J}$ is

$$R_{\ell J} = \frac{\gamma_{\lambda \ell J}^2}{E_{\lambda \ell J} - E}, \quad (B.7)$$

where $\gamma_{\lambda \ell J}^2$ is the reduced width of the resonance, and $E_{\lambda \ell J}$ is the associated energy eigenvalue.

For $\ell = 1$, the explicit forms of F_ℓ and G_ℓ are

$$F_1 = \frac{\sin \rho - \rho \cos \rho}{\rho} \quad (B.8)$$

and

$$G_1 = \frac{\cos \rho + \rho \sin \rho}{\rho}. \quad (B.9)$$

Using Eqs. B.5-B.9 in conjunction with Eq. B.1 and the boundary condition $B_{\ell J} = 0$ results in

$$\delta_{1J} + \phi_1 = \tan^{-1} \left(\frac{\alpha \rho^3}{1 + \beta \rho^2 + (\alpha + \beta - 1) \rho^4} \right) \approx \frac{\alpha \rho^3}{1 + \beta \rho^2 + (\alpha + \beta - 1) \rho^4}, \quad (B.10)$$

where

$$\alpha = \frac{\gamma_{\lambda 1 J}^2}{E_{\lambda 1 J} + \gamma_{\lambda 1 J}^2} \quad (B.11)$$

and

$$\beta = \frac{E_{\lambda 1 J} - \frac{A}{A+1} \frac{1}{C^2}}{E_{\lambda 1 J} + \gamma_{\lambda 1 J}^2} \quad (B.12)$$

The parameters α and β were determined by fitting the low-energy p-wave phase shifts of Table I of Ref. 1 using Eq. B.10 and $a = 3.0 \times 10^{-13}$ cm. The phase shifts at 0.3 and

0.4 MeV were fit for $J = 3/2$; those at 0.4 and 0.6 MeV were fit for $J = 1/2$. This procedure gives two linear equations in α and β for both $J = 3/2$ and $J = 1/2$. Table II presents the solutions of the equations.

TABLE II. Low-energy Parameters for p-wave Phase Shifts

J	α	β
$3/2$	2.528	-2.534
$1/2$	0.857	-0.204

APPENDIX C

Equations for Calculation of $\bar{\mu}_L$, ξ , and γ

Equation 1 gives the angular distribution in the center-of-mass system for elastic scattering. The equations for $\bar{\mu}_L$, ξ , and γ , in terms of the f_ℓ expansion coefficients, derived in Refs. 3-5 and used in FORTRAN program MUXIGA, are given below.

The average cosine of the angle of elastic scattering in the laboratory system is

$$\bar{\mu}_L = \sum_{\ell=0}^{\infty} T_{1\ell} f_\ell, \quad (C.1)$$

where

$$T_{1\ell} = \frac{\ell}{2\ell - 1} \left(\frac{-1}{A}\right)^{\ell-1} - \frac{\ell + 2}{2\ell + 3} \left(\frac{-1}{A}\right)^{\ell+1}, \quad (C.2)$$

A being the ratio of the mass of the scatterer to the neutron mass.

The average logarithmic energy decrement in elastic scattering is

$$\xi = \sum_{\ell=0}^{\infty} \sum_{p=\ell+\delta_{\ell_0}}^{\infty} (1-\alpha)^p A_{p\ell} f_\ell, \quad (C.3)$$

where

$$A_{p\ell} = \frac{(-1)^\ell (2\ell + 1)}{p(p-\ell)! (p+\ell+1)!} p! p!, \quad (C.4)$$

and

$$\alpha = \left(\frac{A - 1}{A + 1}\right)^2.$$

The Greuling-Goertzel parameter, γ , is given by

$$2\xi\gamma = \sum_{\ell=0}^{\infty} \sum_{p=\ell+\delta_{\ell_1+\frac{1}{2}\delta_{\ell_0}}}^{\infty} (1-\alpha)^p A_{p\ell} C_p f_\ell, \quad (C.5)$$

where

$$C_p = 2 \sum_{q=1}^{p-1} \frac{1}{q}. \quad (C.6)$$

The δ_{ℓ_0} and δ_{ℓ_1} in Eqs. C.3 and C.5 are Kronecker δ symbols; that is, $\delta_{\ell\ell'} = 0$ for $\ell \neq \ell'$ and $\delta_{\ell\ell'} = 1$ for $\ell = \ell'$.

APPENDIX D

ENDF/B Listing of Helium

2000.0	3.96822	0	0	0	01088	1451	1
0.0	0.0	0	0	74	01088	1451	2
HELIUM CROSS SECTIONS MATERIAL 1088					1088	1451	3
COMPILED BY ED PENNINGTON, ARGONNE NATIONAL LAB., IN JUNE 1968,					1088	1451	4
NATURAL HELIUM CONSISTS OF 0.00013 PER CENT HE-3 AND 99.99987 PER CENT HE-4,					1088	1451	5
BECAUSE OF THE LOW ABUNDANCE OF HE-3, ONLY ITS (N,P) CROSS SECTION, WHICH IS VERY LARGE AT LOW ENERGIES, NEED BE CONSIDERED.					1088	1451	6
ELASTIC SCATTERING IS THE ONLY POSSIBLE REACTION FOR NEUTRONS INCIDENT ON HE-4 AT ENERGIES BELOW 15 MEV. THUS THE ELASTIC SCATTERING CROSS SECTION AND VALUES OF MU BAR(LAB), XI, AND GAMMA ARE GIVEN IN FILE 3, AND ELASTIC SCATTERING LEGENDRE COEFFICIENTS ARE GIVEN IN FILE 4. PARAMETERS FOR A FREE GAS THERMAL SCATTERING LAW ARE IN FILE 7.					1088	1451	7
THE ELASTIC SCATTERING CROSS SECTION AND THE LEGENDRE EXPANSION COEFFICIENTS WERE CALCULATED FROM S-, P-, AND D-WAVE PHASE SHIFTS USING A FORTRAN PROGRAM WRITTEN FOR THE PURPOSE, THE PHASE SHIFTS WERE READ FROM SMOOTH CURVES BASED ON TABLE I OF REF.1, AT ENERGIES BELOW THE 300 KEV, LOWER LIMIT OF TABLE I, EACH OF THE TWO P-WAVE PHASE SHIFTS WAS OBTAINED BY ASSUMING A FUNCTIONAL FORM BASED ON THE LOW ENERGY LIMIT FOR A SINGLE P-WAVE RESONANCE, WITH PARAMETERS DETERMINED FROM FITTING THE LOW ENERGY PHASE SHIFTS OF TABLE I, THE S-WAVE PHASE SHIFT BELOW 300 KEV, WAS CALCULATED USING HARD SPHERE SCATTERING AND A NUCLEAR RADIUS, A = 2.4 FERMI, THIS YIELDS THE THERMAL SCATTERING CROSS SECTION = 4.*PI*A**2*0.7238 BARNS IN AGREEMENT WITH THE EXPERIMENTAL VALUE OF 0.73** 0.05 BARNS (REF.2). THE LOW ENERGY S-WAVE PHASE SHIFTS OF REF.1 ARE CONSISTENT WITH A NUCLEAR RADIUS OF ABOUT 2.48 FERMI, AND SO WOULD YIELD A SOMEWHAT HIGH THERMAL CROSS SECTION. VALUES OF MU BAR(LAB), XI, AND GAMMA WERE CALCULATED FROM THE LEGENDRE COEFFICIENTS USING A FORTRAN PROGRAM, MUXIGA. THIS PROGRAM USES THE EQUATIONS OF REF.3-5.					1088	1451	14
AN ELASTIC SCATTERING TRANSFORMATION MATRIX FROM THE CENTER-OF-MASS TO THE LABORATORY SYSTEM WAS COMPUTED USING CHAD (REF.6), THE (N,P) CROSS SECTION FOR HE-3 IS THAT RECOMMENDED IN THE EVALUATION OF HE-3 BY J. ALS-NIELSEN GIVEN IN REF.7. EXTENSION FROM 10 TO 15 MEV, WAS MADE USING LINEAR EXTRAPOLATION ON A LOG SIGMA* LOG E SCALE.					1088	1451	15
THE TOTAL CROSS SECTION IS THE SUM OF THE ELASTIC SCATTERING AND (N,P) CROSS SECTIONS.					1088	1451	16
COMMENTS.					1088	1451	17
THE PHASE SHIFTS OF REF.1 ARE OPTICAL MODEL PHASE SHIFTS CHOSEN TO FIT BOTH ANGULAR DISTRIBUTION AND POLARIZATION DATA AT MANY ENERGIES. THE TOTAL SCATTERING CROSS SECTION IS ALSO FIT WITHIN THE SCATTER OF THE EXPERIMENTAL POINTS. ANOTHER RECENT SET OF PHASE SHIFTS (REF.8) IS NOT VERY DIFFERENT FROM THOSE USED HERE, AND COULD ALSO HAVE BEEN USED IN THE PRESENT WORK. THERE SHOULD BE NO SERIOUS DISAGREEMENTS WITH EXPERIMENTAL VALUES IN THE HE-4 DATA CALCULATED FROM THE PHASE SHIFTS, AS DISCUSSED IN REF.7, THE HE-3 (N,P) CROSS SECTION IS RATHER WELL KNOWN, PROBABLY MORE ERROR IS INTRODUCED INTO THE (N,P) CROSS SECTION FOR NATURAL HELIUM BY THE UNCERTAINTY IN THE HE-3 ISOTOPIC ABUNDANCE THAN BY THE UNCERTAINTY IN THE HE-3 (N,P) CROSS SECTION ITSELF.					1088	1451	18
PREVIOUS EVALUATIONS OF HELIUM FOR REACTOR CALCULATIONS INCLUDE THOSE OF J.J.SCHMIDT (REF.9) AND B.R.S.BUCKINGHAM ET AL (REF.10).					1088	1451	19
SCHMIDT'S EVALUATION INCLUDES THE (N,P) CROSS SECTION FOR HE-3, AND SIGMA ELASTIC, MU BAR(LAB), AND A SET OF PHASE SHIFTS FOR HE-4. BUCKINGHAM ET AL GIVE SEPARATE EVALUATIONS FOR HE-3 AND HE-4, FOR HE-3 ELASTIC, (N,P), (N,D), AND (N,2N) CROSS SECTIONS ARE GIVEN, AS WELL AS ELASTIC ANGULAR DISTRIBUTIONS. THE HE-4 EVALUATION GIVES THE ELASTIC CROSS SECTION AND ANGULAR					1088	1451	20

DISTRIBUTIONS,

THE PRESENT COMPILATION WILL BE DESCRIBED IN DETAIL IN REF.11.

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2000.0	3.96822	0	0	0	01088 3	1	79
0.0	0.0	0	0	1	1421088 3	1	80
142	3						
					1088 1		81
1.00E+05	1.0721E+00	2.00E-05	0.9701E+00	3.00E-05	0.9249E+001088 3	1	82
5.00E-05	0.8796E+00	7.00E-05	0.8555E+00	1.00E-04	0.8339E+001088 3	1	83
2.00E-04	0.8017E+00	3.00E-04	0.7874E+00	5.00E-04	0.7731E+001088 3	1	84
7.00E-04	0.7654E+00	1.00E-03	0.7586E+00	2.00E-03	0.7484E+001088 3	1	85
3.00E-03	0.7439E+00	5.00E-03	0.7394E+00	7.00E-03	0.7370E+001088 3	1	86
1.00E-02	0.7348E+00	2.00E-02	0.7316E+00	2.53E-02	0.7307E+001088 3	1	87
3.00E-02	0.7302E+00	5.00E-02	0.7287E+00	7.00E-02	0.7280E+001088 3	1	88
1.00E-01	0.7273E+00	2.00E-01	0.7263E+00	3.00E-01	0.7258E+001088 3	1	89
5.00E-01	0.7254E+00	7.00E-01	0.7251E+00	1.00E+00	0.7249E+001088 3	1	90
2.00E+00	0.7246E+00	3.00E+00	0.7244E+00	5.00E+00	0.7243E+001088 3	1	91
7.00E+00	0.7242E+00	1.00E+01	0.7241E+00	2.00E+01	0.7240E+001088 3	1	92
3.00E+01	0.7240E+00	5.00E+01	0.7240E+00	7.00E+01	0.7239E+001088 3	1	93
1.00E+02	0.7239E+00	2.00E+02	0.7239E+00	3.00E+02	0.7238E+001088 3	1	94
5.00E+02	0.7238E+00	7.00E+02	0.7238E+00	1.00E+03	0.7238E+001088 3	1	95
2.00E+03	0.7237E+00	3.00E+03	0.7237E+00	5.00E+03	0.7236E+001088 3	1	96
7.00E+03	0.7236E+00	1.00E+04	0.7235E+00	2.00E+04	0.7235E+001088 3	1	97
3.00E+04	0.7234E+00	5.00E+04	0.7241E+00	7.00E+04	0.7256E+001088 3	1	98
1.00E+05	0.7297E+00	1.50E+05	0.7425E+00	2.00E+05	0.7644E+001088 3	1	99
2.50E+05	0.7977E+00	3.00E+05	0.8978E+00	3.50E+05	0.9654E+001088 3	1	100
4.00E+05	0.1056E+01	4.50E+05	0.1197E+01	5.00E+05	0.1373E+011088 3	1	101
5.50E+05	0.1581E+01	6.00E+05	0.1826E+01	6.50E+05	0.2168E+011088 3	1	102
7.00E+05	0.2576E+01	7.50E+05	0.3076E+01	8.00E+05	0.3750E+011088 3	1	103
8.50E+05	0.4380E+01	9.00E+05	0.5114E+01	9.50E+05	0.5844E+011088 3	1	104
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1.13E+06	0.7424E+01	1.15E+06	0.7456E+01	1.17E+06	0.7461E+011088 3	1	106
1.20E+06	0.7432E+01	1.25E+06	0.7278E+01	1.30E+06	0.7040E+011088 3	1	107
1.35E+06	0.6757E+01	1.40E+06	0.6453E+01	1.45E+06	0.6157E+011088 3	1	108
1.50E+06	0.5875E+01	1.60E+06	0.5333E+01	1.70E+06	0.4882E+011088 3	1	109
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2.10E+06	0.3734E+01	2.20E+06	0.3551E+01	2.30E+06	0.3399E+011088 3	1	111
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2.70E+06	0.2988E+01	2.80E+06	0.2911E+01	2.90E+06	0.2849E+011088 3	1	113
3.00E+06	0.2789E+01	3.20E+06	0.2684E+01	3.40E+06	0.2595E+011088 3	1	114
3.60E+06	0.2524E+01	3.80E+06	0.2469E+01	4.00E+06	0.2424E+011088 3	1	115
4.20E+06	0.2373E+01	4.40E+06	0.2325E+01	4.60E+06	0.2278E+011088 3	1	116
4.80E+06	0.2232E+01	5.00E+06	0.2190E+01	5.20E+06	0.2149E+011088 3	1	117
5.40E+06	0.2107E+01	5.60E+06	0.2070E+01	5.80E+06	0.2031E+011088 3	1	118
6.00E+06	0.1994E+01	6.20E+06	0.1957E+01	6.40E+06	0.1921E+011088 3	1	119
6.60E+06	0.1887E+01	6.80E+06	0.1853E+01	7.00E+06	0.1821E+011088 3	1	120
7.20E+06	0.1787E+01	7.40E+06	0.1756E+01	7.60E+06	0.1726E+011088 3	1	121
7.80E+06	0.1696E+01	8.00E+06	0.1667E+01	8.20E+06	0.1640E+011088 3	1	122
8.40E+06	0.1612E+01	8.60E+06	0.1585E+01	8.80E+06	0.1559E+011088 3	1	123
9.00E+06	0.1534E+01	9.20E+06	0.1509E+01	9.40E+06	0.1485E+011088 3	1	124

9.60E+06	0.1462E+01	9.80E+06	0.1439E+01	1.00E+07	0.1418E+01	1.088	3	1	125
1.05E+07	0.1345E+01	1.10E+07	0.1315E+01	1.15E+07	0.1268E+01	1.088	3	1	126
1.20E+07	0.1224E+01	1.25E+07	0.1180E+01	1.30E+07	0.1140E+01	1.088	3	1	127
1.35E+07	0.1102E+01	1.40E+07	0.1066E+01	1.45E+07	0.1032E+01	1.088	3	1	128
1.50E+07	0.9992E+00					1.088	3	1	129
0.0	0.0	0	0	0	0	0.088	3	0	130
2000.0	3.96822	0	0	0	0	0.088	3	2	131
0.0	0.0	0	0	1	1.421088	3	2	132	
142	3					1.088	3	2	133
1.00E-05	0.7238E+00	2.00E-05	0.7238E+00	3.00E-05	0.7238E+00	1.088	3	2	134
5.00E-05	0.7238E+00	7.00E-05	0.7238E+00	1.00E-04	0.7238E+00	1.088	3	2	135
2.00E-04	0.7238E+00	3.00E-04	0.7238E+00	5.00E-04	0.7238E+00	1.088	3	2	136
7.00E-04	0.7238E+00	1.00E-03	0.7238E+00	2.00E-03	0.7238E+00	1.088	3	2	137
3.00E-03	0.7238E+00	5.00E-03	0.7238E+00	7.00E-03	0.7238E+00	1.088	3	2	138
1.00E-02	0.7238E+00	2.00E-02	0.7238E+00	2.53E-02	0.7238E+00	1.088	3	2	139
3.00E-02	0.7238E+00	5.00E-02	0.7238E+00	7.00E-02	0.7238E+00	1.088	3	2	140
1.00E-01	0.7238E+00	2.00E-01	0.7238E+00	3.00E-01	0.7238E+00	1.088	3	2	141
5.00E-01	0.7238E+00	7.00E-01	0.7238E+00	1.00E+00	0.7238E+00	1.088	3	2	142
2.00E+00	0.7238E+00	3.00E+00	0.7238E+00	5.00E+00	0.7238E+00	1.088	3	2	143
7.00E+00	0.7238E+00	1.00E+01	0.7238E+00	2.00E+01	0.7238E+00	1.088	3	2	144
3.00E+01	0.7238E+00	5.00E+01	0.7238E+00	7.00E+01	0.7238E+00	1.088	3	2	145
1.00E+02	0.7238E+00	2.00E+02	0.7238E+00	3.00E+02	0.7238E+00	1.088	3	2	146
5.00E+02	0.7238E+00	7.00E+02	0.7238E+00	1.00E+03	0.7238E+00	1.088	3	2	147
2.00E+03	0.7237E+00	3.00E+03	0.7237E+00	5.00E+03	0.7236E+00	1.088	3	2	148
7.00E+03	0.7236E+00	1.00E+04	0.7235E+00	2.00E+04	0.7235E+00	1.088	3	2	149
3.00E+04	0.7234E+00	5.00E+04	0.7234E+00	7.00E+04	0.7256E+00	1.088	3	2	150
1.00E+05	0.7297E+00	1.50E+05	0.7425E+00	2.00E+05	0.7644E+00	1.088	3	2	151
2.50E+05	0.7977E+00	3.00E+05	0.8978E+00	3.50E+05	0.9654E+00	1.088	3	2	152
4.00E+05	0.1056E+01	4.50E+05	0.1197E+01	5.00E+05	0.1373E+01	1.088	3	2	153
5.50E+05	0.1581E+01	6.00E+05	0.1826E+01	6.50E+05	0.2168E+01	1.088	3	2	154
7.00E+05	0.2576E+01	7.50E+05	0.3076E+01	8.00E+05	0.3750E+01	1.088	3	2	155
8.50E+05	0.4380E+01	9.00E+05	0.5114E+01	9.50E+05	0.5844E+01	1.088	3	2	156
1.00E+06	0.6525E+01	1.05E+06	0.7038E+01	1.10E+06	0.7321E+01	1.088	3	2	157
1.13E+06	0.7424E+01	1.15E+06	0.7456E+01	1.17E+06	0.7461E+01	1.088	3	2	158
1.20E+06	0.7432E+01	1.25E+06	0.7278E+01	1.30E+06	0.7040E+01	1.088	3	2	159
1.35E+06	0.6757E+01	1.40E+06	0.6453E+01	1.45E+06	0.6157E+01	1.088	3	2	160
1.50E+06	0.5875E+01	1.60E+06	0.5333E+01	1.70E+06	0.4882E+01	1.088	3	2	161
1.80E+06	0.4509E+01	1.90E+06	0.4211E+01	2.00E+06	0.3944E+01	1.088	3	2	162
2.10E+06	0.3734E+01	2.20E+06	0.3551E+01	2.30E+06	0.3399E+01	1.088	3	2	163
2.40E+06	0.3273E+01	2.50E+06	0.3169E+01	2.60E+06	0.3071E+01	1.088	3	2	164
2.70E+06	0.2988E+01	2.80E+06	0.2911E+01	2.90E+06	0.2849E+01	1.088	3	2	165
3.00E+06	0.2789E+01	3.20E+06	0.2628E+01	3.40E+06	0.2595E+01	1.088	3	2	166
3.60E+06	0.2524E+01	3.80E+06	0.2469E+01	4.00E+06	0.2424E+01	1.088	3	2	167
4.20E+06	0.2373E+01	4.40E+06	0.2325E+01	4.60E+06	0.2278E+01	1.088	3	2	168
4.80E+06	0.2232E+01	5.00E+06	0.2190E+01	5.20E+06	0.2149E+01	1.088	3	2	169
5.40E+06	0.2107E+01	5.60E+06	0.2070E+01	5.80E+06	0.2031E+01	1.088	3	2	170
6.00E+06	0.1994E+01	6.20E+06	0.1957E+01	6.40E+06	0.1921E+01	1.088	3	2	171
6.60E+06	0.1887E+01	6.80E+06	0.1853E+01	7.00E+06	0.1821E+01	1.088	3	2	172
7.20E+06	0.1787E+01	7.40E+06	0.1756E+01	7.60E+06	0.1726E+01	1.088	3	2	173
7.80E+06	0.1696E+01	8.00E+06	0.1667E+01	8.20E+06	0.1640E+01	1.088	3	2	174
8.40E+06	0.1612E+01	8.60E+06	0.1585E+01	8.80E+06	0.1559E+01	1.088	3	2	175
9.00E+06	0.1534E+01	9.20E+06	0.1509E+01	9.40E+06	0.1485E+01	1.088	3	2	176
9.60E+06	0.1462E+01	9.80E+06	0.1439E+01	1.00E+07	0.1418E+01	1.088	3	2	177
1.05E+07	0.1365E+01	1.10E+07	0.1315E+01	1.15E+07	0.1268E+01	1.088	3	2	178
1.20E+07	0.1224E+01	1.25E+07	0.1180E+01	1.30E+07	0.1140E+01	1.088	3	2	179
1.35E+07	0.1102E+01	1.40E+07	0.1066E+01	1.45E+07	0.1032E+01	1.088	3	2	180
1.50E+07	0.9992E+00					1.088	3	2	181
0.0	0.0	0	0	0	0	0.088	3	0	182
2000.0	3.96822	0	0	0	0	0.088	3103	183	
0.0	0.7644E+06	0	0	1	1.421088	3103	184		
142	5					1.088	3103	185	
1.00E-05	0.34483E+00	2.00E-05	0.2463E+00	3.00E-05	0.2011E+00	1.088	3103	186	

5.00E-05	0.1558E+00	7.00E-05	0.1317E+00	1.00E-04	0.1101E+00	3103	187
2.00E-04	0.7789E-01	3.00E-04	0.6359E-01	5.00E-04	0.4926E-01	3103	188
7.00E-04	0.4163E-01	1.00E-03	0.3483E-01	2.00E-03	0.2463E-01	3103	189
3.00E-03	0.2011E-01	5.00E-03	0.1558E-01	7.00E-03	0.1316E-01	3103	190
1.00E-02	0.1101E-01	2.00E-02	0.7787E-02	2.53E-02	0.6924E-02	3103	191
3.00E-02	0.6358E-02	5.00E-02	0.4924E-02	7.00E-02	0.4161E-02	3103	192
1.00F+01	0.3481E-02	2.00E-01	0.2461E-02	3.00E-01	0.2010E-02	3103	193
5.00F-01	0.1556E-02	7.00E-01	0.1316E-02	1.00E+00	0.1100E-02	3103	194
2.00E+00	0.7773E-03	3.00E+00	0.6344E-03	5.00E+00	0.4910E-03	3103	195
7.00F+00	0.4147E-03	1.00E+01	0.3467E-03	2.00E+01	0.2448E-03	3103	196
3.00E+01	0.1996E-03	5.00E+01	0.1542E-03	7.00E+01	0.1301E-03	3103	197
1.00E+02	0.1086E-03	2.00E+02	0.7631E-04	3.00E+02	0.6201E-04	3103	198
5.00E+02	0.4771E-04	7.00E+02	0.4004E-04	1.00E+03	0.3328E-04	3103	199
2.00E+03	0.2314E-04	3.00E+03	0.1859E-04	5.00E+03	0.1417E-04	3103	200
7.00E+03	0.1183E-04	1.00E+04	0.9880E-05	2.00E+04	0.6825E-05	3103	201
3.00E+04	0.5330E-05	5.00E+04	0.3900E-05	7.00E+04	0.3185E-05	3103	202
1.00E+05	0.2574E-05	1.50E+05	0.2067E-05	2.00E+05	0.1775E-05	3103	203
2.50E+05	0.1560E-05	3.00E+05	0.1430E-05	3.50E+05	0.1313E-05	3103	204
4.00E+05	0.1248E-05	4.50E+05	0.1196E-05	5.00E+05	0.1183E-05	3103	205
5.50E+05	0.1170E-05	6.00E+05	0.1157E-05	6.50E+05	0.1157E-05	3103	206
7.00E+05	0.1157E-05	7.50E+05	0.1157E-05	8.00E+05	0.1157E-05	3103	207
8.50E+05	0.1157E-05	9.00E+05	0.1157E-05	9.50E+05	0.1157E-05	3103	208
1.00E+06	0.1157E-05	1.05E+06	0.1157E-05	1.10E+06	0.1157E-05	3103	209
1.13E+06	0.1157E-05	1.15E+06	0.1157E-05	1.17E+06	0.1157E-05	3103	210
1.20E+06	0.1157E-05	1.25E+06	0.1157E-05	1.30E+06	0.1157E-05	3103	211
1.35E+06	0.1157E-05	1.40E+06	0.1157E-05	1.45E+06	0.1157E-05	3103	212
1.50E+06	0.1157E-05	1.60E+06	0.1151E-05	1.70E+06	0.1144E-06	3103	213
1.80E+06	0.1138E-05	1.90E+06	0.1118E-05	2.00E+06	0.1092E-06	3103	214
2.10F+06	0.1066E-05	2.20E+06	0.1040E-05	2.30E+06	0.1014E-05	3103	215
2.40E+06	0.9880E-06	2.50E+06	0.9520E-06	2.60E+06	0.9360E-06	3103	216
2.70E+06	0.9100E-06	2.80E+06	0.8840E-06	2.90E+06	0.8580E-06	3103	217
3.00E+06	0.8320E-06	3.20E+06	0.7800E-06	3.40E+06	0.7280E-06	3103	218
3.60E+06	0.6825E-06	3.80E+06	0.6433E-06	4.00E+06	0.6110E-06	3103	219
4.20E+06	0.5785E-06	4.40E+06	0.5460E-06	4.60E+06	0.5265E-06	3103	220
4.80E+06	0.5005E-06	5.00E+06	0.4810E-06	5.20E+06	0.4615E-06	3103	221
5.40E+06	0.4420E-06	5.60E+06	0.4225E-06	5.80E+06	0.4095E-06	3103	222
6.00E+06	0.3956E-06	6.20E+06	0.3770E-06	6.40E+06	0.3640E-06	3103	223
6.60E+06	0.3510E-06	6.80E+06	0.3406E-06	7.00E+06	0.3315E-06	3103	224
7.20E+06	0.3211E-06	7.40E+06	0.3133E-06	7.60E+06	0.3055E-06	3103	225
7.80E+06	0.2977E-06	8.00E+06	0.2912E-06	8.20E+06	0.2847E-06	3103	226
8.40E+06	0.2749E-06	8.60E+06	0.2717E-06	8.80E+06	0.2665E-06	3103	227
9.00E+06	0.2600E-06	9.20E+06	0.2548E-06	9.40E+06	0.2483E-06	3103	228
9.60E+06	0.2431E-06	9.80E+06	0.2379E-06	1.00E+07	0.2340E-06	3103	229
1.05E+07	0.2236E-06	1.10E+07	0.2145E-06	1.15E+07	0.2054E-06	3103	230
1.20E+07	0.1976E-06	1.25E+07	0.1898E-06	1.30E+07	0.1833E-06	3103	231
1.35E+07	0.1768E-06	1.40E+07	0.1703E-06	1.45E+07	0.1651E-06	3103	232
1.50E+07	0.1599E-06	0.0	0	0	0	1088	3103
2000.0	3.96822	0	0	0	0	1088	3
0.0	0.0	0	0	0	0	1088	3251
108	3	0	0	1	1	1081088	3251
1.00E-05	0.1680E+00	2.53E-02	0.1680E+00	1.00E+02	0.1679E+00	3251	238
2.00F+02	0.1678E+00	3.00E+02	0.1677E+00	5.00E+02	0.1675E+00	3251	239
7.00E+02	0.1672E+00	1.00E+03	0.1669E+00	2.00E+03	0.1658E+00	3251	240
3.00E+03	0.1647E+00	5.00E+03	0.1625E+00	7.00E+03	0.1603E+00	3251	241
1.00E+04	0.1570E+00	2.00E+04	0.1458E+00	3.00E+04	0.1344E+00	3251	242
5.00E+04	0.1116E+00	7.00E+04	0.8798E-01	1.00E+05	0.5208E+00	3251	243
1.50E+05	0.9740E-02	2.00E+05	0.7161E-01	2.50E+05	0.1304E+00	3251	244
3.00E+05	0.1739E+00	3.50E+05	0.2185E+00	4.00E+05	0.2518E+00	3251	245
4.50F+05	0.2702E+00	5.00E+05	0.2746E+00	5.50E+05	0.2658E+00	3251	246
6.00E+05	0.2457E+00	6.50E+05	0.2143E+00	7.00E+05	0.1776E+00	3251	247
7.50F+05	0.1357E+00	8.00E+05	0.8704E-01	8.50E+05	0.4441E-01	3251	248

9.00E+05	0.5200E-04	9.50E+05	0.4267E-01	1.00E+06	0.8390E-01	1.088	3251	249
1.05E+06	0.1213E+00	1.10E+06	0.1529E+00	1.13E+06	0.1720E+00	1.088	3251	250
1.15E+06	0.1836E+00	1.17E+06	0.1948E+00	1.20E+06	0.2115E+00	1.088	3251	251
1.25E+06	0.2370E+00	1.30E+06	0.2585E+00	1.35E+06	0.2784E+00	1.088	3251	252
1.40E+06	0.2967E+00	1.45E+06	0.3122E+00	1.50E+06	0.3258E+00	1.088	3251	253
1.60E+06	0.3520E+00	1.70E+06	0.3718E+00	1.80E+06	0.3878E+00	1.088	3251	254
1.90E+06	0.4002E+00	2.00E+06	0.4104E+00	2.10E+06	0.4185E+00	1.088	3251	255
2.20E+06	0.4255E+00	2.30E+06	0.4315E+00	2.40E+06	0.4363E+00	1.088	3251	256
2.50E+06	0.4402E+00	2.60E+06	0.4434E+00	2.70E+06	0.4460E+00	1.088	3251	257
2.80E+06	0.4489E+00	2.90E+06	0.4512E+00	3.00E+06	0.4530E+00	1.088	3251	258
3.20E+06	0.4585E+00	3.40E+06	0.4634E+00	3.60E+06	0.4675E+00	1.088	3251	259
3.80E+06	0.4710E+00	4.00E+06	0.4749E+00	4.20E+06	0.4795E+00	1.088	3251	260
4.40E+06	0.4845E+00	4.60E+06	0.4892E+00	4.80E+06	0.4940E+00	1.088	3251	261
5.00E+06	0.4987E+00	5.20E+06	0.5035E+00	5.40E+06	0.5077E+00	1.088	3251	262
5.60E+06	0.5117E+00	5.80E+06	0.5153E+00	6.00E+06	0.5195E+00	1.088	3251	263
6.20E+06	0.5231E+00	6.40E+06	0.5264E+00	6.60E+06	0.5301E+00	1.088	3251	264
6.80E+06	0.5333E+00	7.00E+06	0.5365E+00	7.20E+06	0.5390E+00	1.088	3251	265
7.40E+06	0.5418E+00	7.60E+06	0.5443E+00	7.80E+06	0.5470E+00	1.088	3251	266
8.00E+06	0.5492E+00	8.20E+06	0.5513E+00	8.40E+06	0.5530E+00	1.088	3251	267
8.60E+06	0.5549E+00	8.80E+06	0.5565E+00	9.00E+06	0.5584E+00	1.088	3251	268
9.20E+06	0.5630E+00	9.40E+06	0.5618E+00	9.50E+06	0.5635E+00	1.088	3251	269
9.80E+06	0.5649E+00	1.00E+07	0.5666E+00	1.05E+07	0.5696E+00	1.088	3251	270
1.10E+07	0.5728E+00	1.15E+07	0.5756E+00	1.20E+07	0.5783E+00	1.088	3251	271
1.25E+07	0.5804E+00	1.30E+07	0.5828E+00	1.35E+07	0.5850E+00	1.088	3251	272
1.40E+07	0.5872E+00	1.45E+07	0.5894E+00	1.50E+07	0.5917E+00	1.088	3251	273
0.0	0.0	0	0	0	0	1.088	3	0
2.000.0	3.96822	0	0	0	0	1.088	3252	275
0.0	0.0	0	0	1	1.088	1.088	3252	276
108	3				1.088	3252	277	
1.00E-05	0.4282E+00	2.53E-02	0.4282E+00	1.00E+02	0.4282E+00	1.088	3252	278
2.00E+02	0.4283E+00	3.00E+02	0.4284E+00	5.00E+02	0.4285E+00	1.088	3252	279
7.00E+02	0.4286E+00	1.00E+03	0.4287E+00	2.00E+03	0.4293E+00	1.088	3252	280
3.00E+03	0.4299E+00	5.00E+03	0.4310E+00	7.00E+03	0.4322E+00	1.088	3252	281
1.00E+04	0.4339E+00	2.00E+04	0.4397E+00	3.00E+04	0.4455E+00	1.088	3252	282
5.00E+04	0.4574E+00	7.00E+04	0.4696E+00	1.00E+05	0.4881E+00	1.088	3252	283
1.50E+05	0.5201E+00	2.00E+05	0.5521E+00	2.50E+05	0.5826E+00	1.088	3252	284
3.00E+05	0.6051E+00	3.50E+05	0.6281E+00	4.00E+05	0.6454E+00	1.088	3252	285
4.50E+05	0.6550E+00	5.00E+05	0.6572E+00	5.50E+05	0.6527E+00	1.088	3252	286
6.00E+05	0.6423E+00	6.50E+05	0.6261E+00	7.00E+05	0.6071E+00	1.088	3252	287
7.50E+05	0.5855E+00	8.00E+05	0.5603F+00	8.50E+05	0.5382E+00	1.088	3252	288
9.00E+05	0.5152E+00	9.50E+05	0.4932E+00	1.00E+06	0.4719E+00	1.088	3252	289
1.05E+06	0.4525E+00	1.10E+06	0.4362E+00	1.13E+06	0.4263E+00	1.088	3252	290
1.15E+06	0.4203E+00	1.17E+06	0.4145E+00	1.20E+06	0.4059E+00	1.088	3252	291
1.25E+06	0.3927E+00	1.30E+06	0.3816E+00	1.35E+06	0.3712E+00	1.088	3252	292
1.40E+06	0.3618E+00	1.45E+06	0.3538E+00	1.50E+06	0.3467E+00	1.088	3252	293
1.60E+06	0.3331E+00	1.70E+06	0.3229E+00	1.80E+06	0.3146E+00	1.088	3252	294
1.90E+06	0.3028E+00	2.00E+06	0.3029E+00	2.10E+06	0.2988E+00	1.088	3252	295
2.20E+06	0.2951E+00	2.30E+06	0.2920E+00	2.40E+06	0.2895E+00	1.088	3252	296
2.50E+06	0.2875E+00	2.60E+06	0.2858E+00	2.70E+06	0.2845E+00	1.088	3252	297
2.80E+06	0.2830E+00	2.90E+06	0.2818E+00	3.00E+06	0.2809E+00	1.088	3252	298
3.20E+06	0.2781E+00	3.40E+06	0.2755E+00	3.50E+06	0.2734E+00	1.088	3252	299
3.80E+06	0.2716E+00	4.00E+06	0.2696E+00	4.20E+06	0.2672E+00	1.088	3252	300
4.40E+06	0.2646E+00	4.60E+06	0.2622E+00	4.80E+06	0.2597E+00	1.088	3252	301
5.00E+06	0.2573E+00	5.20E+06	0.2548E+00	5.40E+06	0.2526E+00	1.088	3252	302
5.60E+06	0.2506E+00	5.80E+06	0.2487E+00	6.00E+06	0.2465E+00	1.088	3252	303
6.20E+06	0.2447E+00	6.40E+06	0.2430E+00	6.50E+06	0.2411E+00	1.088	3252	304
6.80E+06	0.2394E+00	7.00E+06	0.2378E+00	7.20E+06	0.2365E+00	1.088	3252	305
7.40E+06	0.2350E+00	7.60E+06	0.2337E+00	7.80E+06	0.2323E+00	1.088	3252	306
8.00E+06	0.2312E+00	8.20E+06	0.2301E+00	8.40E+06	0.2293E+00	1.088	3252	307
8.60E+06	0.2293E+00	8.80E+06	0.2274E+00	9.00E+06	0.2265E+00	1.088	3252	308
9.20E+06	0.2255E+00	9.40E+06	0.2247E+00	9.50E+06	0.2238E+00	1.088	3252	309
9.80E+06	0.2231E+00	1.00E+07	0.2222E+00	1.05E+07	0.2207E+00	1.088	3252	310

1.10E+07	0.2190E+00	1.15E+07	0.2176E+00	1.20E+07	0.2162E+00	1.088	3	0	311
1.25F+07	0.2151E+00	1.30E+07	0.2139E+00	1.35E+07	0.2127E+00	1.088	3	0	312
1.40E+07	0.2116E+00	1.45E+07	0.2104E+00	1.50E+07	0.2093E+00	1.088	3	0	313
0.0	0.0	0	0	0	0	0	1	1	314
2000.0	3.96822	0	0	0	0	0	1	1	315
0.0	0.0	0	0	0	0	0	1	1	316
108	3					1088	3	2	317
1.00E-05	0.3121E+00	2.53E-02	0.3121E+00	1.00E+02	0.3121E+00	1.088	3	2	318
2.00E+02	0.3122E+00	3.00E+02	0.3122E+00	5.00E+02	0.3122E+00	1.088	3	2	319
7.00E+02	0.3122E+00	1.00E+03	0.3123E+00	2.00E+03	0.3125E+00	1.088	3	2	320
3.00E+03	0.3127E+00	5.00E+03	0.3131E+00	7.00E+03	0.3135E+00	1.088	3	2	321
1.00E+04	0.3141E+00	2.00E+04	0.3161E+00	3.00E+04	0.3180E+00	1.088	3	2	322
5.00E+04	0.3219E+00	7.00E+04	0.3257E+00	1.00E+05	0.3312E+00	1.088	3	2	323
1.50E+05	0.3402E+00	2.00E+05	0.3485E+00	2.50E+05	0.3561E+00	1.088	3	2	324
3.00E+05	0.3617E+00	3.50E+05	0.3675E+00	4.00E+05	0.3724E+00	1.088	3	2	325
4.50E+05	0.3762E+00	5.00E+05	0.3790E+00	5.50E+05	0.3808E+00	1.088	3	2	326
6.00E+05	0.3818E+00	6.50E+05	0.3817E+00	7.00E+05	0.3809E+00	1.088	3	2	327
7.50E+05	0.3791E+00	8.00E+05	0.3765E+00	8.50E+05	0.3736E+00	1.088	3	2	328
9.00E+05	0.3702E+00	9.50E+05	0.3662E+00	1.00E+06	0.3618E+00	1.088	3	2	329
1.05E+06	0.3572E+00	1.10E+06	0.3530E+00	1.13E+06	0.3501E+00	1.088	3	2	330
1.15E+06	0.3482E+00	1.17E+06	0.3464E+00	1.20E+06	0.3435E+00	1.088	3	2	331
1.25E+06	0.3388E+00	1.30E+06	0.3344E+00	1.35E+06	0.3300E+00	1.088	3	2	332
1.40E+06	0.3256E+00	1.45E+06	0.3216E+00	1.50E+06	0.3179E+00	1.088	3	2	333
1.60E+06	0.3100E+00	1.70E+06	0.3034E+00	1.80E+06	0.2975E+00	1.088	3	2	334
1.90E+06	0.2925E+00	2.00E+06	0.2879E+00	2.10E+06	0.2842E+00	1.088	3	2	335
2.20E+06	0.2808E+00	2.30E+06	0.2779E+00	2.40E+06	0.2756E+00	1.088	3	2	336
2.50E+06	0.2738E+00	2.60E+06	0.2722E+00	2.70E+06	0.2711E+00	1.088	3	2	337
2.80E+06	0.2699E+00	2.90E+06	0.2691E+00	3.00E+06	0.2687E+00	1.088	3	2	338
3.20E+06	0.2671E+00	3.40E+06	0.2661E+00	3.60E+06	0.2658E+00	1.088	3	2	339
3.80E+06	0.2664E+00	4.00E+06	0.2671E+00	4.20E+06	0.2672E+00	1.088	3	2	340
4.40E+06	0.2671E+00	4.60E+06	0.2671E+00	4.80E+06	0.2670E+00	1.088	3	2	341
5.00E+06	0.2669E+00	5.20E+06	0.2666E+00	5.40E+06	0.2665E+00	1.088	3	2	342
5.60E+06	0.2664E+00	5.80E+06	0.2664E+00	6.00E+06	0.2661E+00	1.088	3	2	343
6.20E+06	0.2659E+00	6.40E+06	0.2656E+00	6.60E+06	0.2652E+00	1.088	3	2	344
6.80E+06	0.2648E+00	7.00E+06	0.2644E+00	7.20E+06	0.2641E+00	1.088	3	2	345
7.40E+06	0.2636E+00	7.60E+06	0.2634E+00	7.80E+06	0.2628E+00	1.088	3	2	346
8.00E+06	0.2624E+00	8.20E+06	0.2622E+00	8.40E+06	0.2620E+00	1.088	3	2	347
8.60E+06	0.2616E+00	8.80E+06	0.2614E+00	9.00E+06	0.2610E+00	1.088	3	2	348
9.20E+06	0.2604E+00	9.40E+06	0.2601F+00	9.60E+06	0.2597E+00	1.088	3	2	349
9.80E+06	0.2594E+00	1.00E+07	0.2590E+00	1.05E+07	0.2583E+00	1.088	3	2	350
1.10E+07	0.2576E+00	1.15E+07	0.2568E+00	1.20E+07	0.2559E+00	1.088	3	2	351
1.25E+07	0.2552E+00	1.30E+07	0.2543E+00	1.35E+07	0.2536E+00	1.088	3	2	352
1.40E+07	0.2529E+00	1.45E+07	0.2522E+00	1.50E+07	0.2515E+00	1.088	3	2	353
0.0	0.0	0	0	0	0	0	1	0	354
0.0	0.0	0	0	0	0	0	1	0	355
2000.0	3.96822	1	1	0	0	0	1	0	356
0.0	3.96822	0	2	100	0	0	1	0	357
1.0000E+00	1.6800E-01	1.2819E-02	0.0-6.6316E-05			0.01088	4	2	358
6.3701E-07	0.0-7.4839E-09	0.0	0.0	9.6190E-01	0.01088	4	2	359	
2.9686E-01	4.3546E-02	3.1027E-03	0.0-1.4971E-05			0.01088	4	2	360
1.3935E-07	0.0	0.0-1.5886E-01	9.0195E-01	4.1067E-01	0.01088	4	2	361	
8.8718E-02	1.1085E-02	7.1925E-04	0.0-3.1239E-06			0.01088	4	2	362
0.0	3.5863E-02-2.7304E-01	8.1319E-01	5.0658E-01	1.4568E-01	0.01088	4	2	363	
2.5474E-02	2.8077E-03	1.6752E-04	0.0	0.0-5.5905E-03	0.01088	4	2	364	
7.8689E-02	-3.6481E-01	7.0133E-01	5.8110E-01	2.1109E-01	4.7327E-02	0.01088	4	2	365
7.1261E-03	7.0975E-04	0.0	2.1026E-03-2.2043E-02	1.2928E-01	0.01088	4	2	366	
-4.3373E-01	5.7179E-01	6.3117E-01	2.8107E-01	7.7131E-02	1.4600E-02	0.01088	4	2	367
0.0-5.1992E-04	6.0621E-03-4.1723E-02	1.8451E-01	4.7795E-01	0.01088	4	2	368		
4.3070E-01	6.54579E-01	3.5139E-01	1.1474E-01	0.0	1.2929E-04	0.01088	4	2	369
-1.6456E-03	1.2786E-02-6.7791E-02	2.4055E-01	4.9642E-01	2.8461E-01	0.01088	4	2	370	
6.5122E-01	4.1767E-01	0.0-3.2264E-05	4.4243E-04	3.7901E-03	0.01088	4	2	371	
2.2986E-02	-9.9735E-02	2.9341E-01	4.8921E-01	1.4018E-01	6.2111E-01	0.01088	4	2	372

0.0	8.0695E-06	-1.1807E-04	1.0976E-03	-7.4043E-03	3.7198E-02	10888	4	2	373
-1.3644E-01	3.3928E-01	-4.5769E-01	3.8210E-03			10888	4	2	374
0.0	0.0	0	0	1		10888	4	2	375
108	3					10888	4	2	376
0.0	1.00E-05	0	0	1		01088	4	2	377
0.0	2.53E-02	0	0	1		01088	4	2	378
0.0						01088	4	2	379
0.0	1.00E+02	0	0	2		01088	4	2	380
-1.130E-04	0.0					01088	4	2	381
0.0	2.00E+02	0	0	2		01088	4	2	382
-2.270E-04	0.0					10888	4	2	384
0.0	3.00E+02	0	0	2		01088	4	2	385
-3.400E-04	0.0					10888	4	2	386
0.0	5.00E+02	0	0	2		01088	4	2	387
-5.670E-04	0.0					10888	4	2	388
0.0	7.00E+02	0	0	2		01088	4	2	389
-7.940E-04	0.0					10888	4	2	390
0.0	1.00E+03	0	0	2		01088	4	2	391
-1.135E-03	0.0					10888	4	2	392
0.0	2.00E+03	0	0	2		01088	4	2	393
-2.272E-03	1.000E-06					10888	4	2	394
0.0	3.00E+03	0	0	2		01088	4	2	395
-3.410E-03	3.000E-06					10888	4	2	396
0.0	5.00E+03	0	0	2		01088	4	2	397
-5.691E-03	9.000E-06					10888	4	2	398
0.0	7.00E+03	0	0	2		01088	4	2	399
-7.979E-03	1.700E-05					10888	4	2	400
0.0	1.00E+04	0	0	2		01088	4	2	401
-1.142E-02	3.500E-05					10888	4	2	402
0.0	2.00E+04	0	0	2		01088	4	2	403
-2.303E-02	1.410E-04					10888	4	2	404
0.0	3.00E+04	0	0	2		01088	4	2	405
-3.483E-02	3.230E-04					10888	4	2	406
0.0	5.00E+04	0	0	2		01088	4	2	407
-5.853E-02	9.160E-04					10888	4	2	408
0.0	7.00E+04	0	0	2		01088	4	2	409
-8.289E-02	1.850E-03					10888	4	2	410
0.0	1.00E+05	0	0	2		01088	4	2	411
-1.199E-01	3.916E-03					10888	4	2	412
0.0	1.50E+05	0	0	2		01088	4	2	413
-1.832E-01	9.426E-03					10888	4	2	414
0.0	2.00E+05	0	0	2		01088	4	2	415
-2.462E-01	1.777E-02					10888	4	2	416
0.0	2.50E+05	0	0	2		01088	4	2	417
-3.054E-01	2.902E-02					10888	4	2	418
0.0	3.00E+05	0	0	2		01088	4	2	419
-3.487E-01	4.071E-02					10888	4	2	420
0.0	3.50E+05	0	0	2		01088	4	2	421
-3.924E-01	5.688E-02					10888	4	2	422
0.0	4.00E+05	0	0	2		01088	4	2	423
-4.240E-01	7.510E-02					10888	4	2	424
0.0	4.50E+05	0	0	2		01088	4	2	425
-4.401E-01	9.412E-02					10888	4	2	426
0.0	5.00E+05	0	0	2		01088	4	2	427
-4.414E-01	1.134E-01					10888	4	2	428
0.0	5.50E+05	0	0	2		01088	4	2	429
-4.291E-01	1.322E-01					10888	4	2	430
0.0	6.00E+05	0	0	2		01088	4	2	431
-4.053E-01	1.499E-01					10888	4	2	432
0.0	6.50E+05	0	0	2		01088	4	2	433
-3.704E-01	1.641E-01					10888	4	2	434

	0.0	7.000E+05	0	0	2				
-3.301E-01	1.769E-01		0			01088	4	2	435
0.0	7.500E+05		0		2	1088	4	2	436
-2.852E-01	1.852E-01		0			01088	4	2	437
0.0	8.000E+05		0		2	1088	4	2	438
-2.333E-01	1.929E-01		0			01088	4	2	439
0.0	8.500E+05		0		2	1088	4	2	440
-1.884E-01	1.967E-01		0			01088	4	2	441
0.0	9.000E+05		0		2	1088	4	2	442
-1.416E-01	1.997E-01		0			01088	4	2	443
0.0	9.500E+05		0		2	1088	4	2	444
-9.717E-02	2.006E-01		0			01088	4	2	445
0.0	1.000E+06		0		2	1088	4	2	446
-5.435E-02	2.003E-01		0			01088	4	2	447
0.0	1.050E+06		0		2	1088	4	2	448
-1.572E-02	1.988E-01		0			01088	4	2	449
0.0	1.100E+06		0		2	1088	4	2	450
1.681E-02	1.970E-01		0			01088	4	2	451
0.0	1.130E+06		0		2	1088	4	2	452
3.635E-02	1.952E-01		0			01088	4	2	453
0.0	1.150E+06		0		2	1088	4	2	454
4.831E-02	1.941E-01		0			01088	4	2	455
0.0	1.170E+06		0		2	1088	4	2	456
5.970E-02	1.929E-01		0			01088	4	2	457
0.0	1.200E+06		0		2	1088	4	2	458
7.677E-02	1.910E-01		0			01088	4	2	459
0.0	1.250E+06		0		2	1088	4	2	460
1.026E-01	1.873E-01		0			01088	4	2	461
0.0	1.300E+06		0		2	1088	4	2	462
1.244E-01	1.837E-01		0			01088	4	2	463
0.0	1.350E+06		0		2	1088	4	2	464
1.445E-01	1.800E-01		0			01088	4	2	465
0.0	1.400E+06		0		2	1088	4	2	466
1.629E-01	1.760E-01		0			01088	4	2	467
0.0	1.450E+06		0		2	1088	4	2	468
1.783E-01	1.723E-01		0			01088	4	2	469
0.0	1.500E+06		0		2	1088	4	2	470
1.919E-01	1.687E-01		0			01088	4	2	471
0.0	1.600E+06		0		2	1088	4	2	472
2.179E-01	1.608E-01		0			01088	4	2	473
0.0	1.700E+06		0		2	1088	4	2	474
2.372E-01	1.538E-01		0			01088	4	2	475
0.0	1.800E+06		0		2	1088	4	2	476
2.528E-01	1.474E-01		0			01088	4	2	477
0.0	1.900E+06		0		2	01088	4	2	478
2.649E-01	1.419E-01		0			1088	4	2	479
0.0	2.000E+06		0		4	01088	4	2	480
2.745E-01	1.357E-01	-2.430E-04	1.0000E-06			1088	4	2	481
0.0	2.100E+06		0		4	01088	4	2	482
2.821E-01	1.314E-01	-2.440E-04	1.0000E-06			1088	4	2	483
0.0	2.200E+06		0		4	01088	4	2	484
2.888E-01	1.272E-01	-2.730E-04	1.0000E-06			1088	4	2	485
0.0	2.300E+06		0		4	01088	4	2	486
2.944E-01	1.239E-01	-3.210E-04	2.0000E-06			1088	4	2	487
0.0	2.400E+06		0		4	01088	4	2	488
2.989E-01	1.213E-01	-3.130E-04	2.0000E-06			1088	4	2	489
0.0	2.500E+06		0		4	01088	4	2	490
3.028E-01	1.195E-01	-3.510E-04	2.0000E-06			01088	4	2	491
0.0	2.600E+06		0		4	1088	4	2	492
3.058E-01	1.180E-01	-3.870E-04	3.0000E-06			01088	4	2	493
0.0	2.700E+06		0		4	1088	4	2	494
3.084E-01	1.171E-01	-3.940E-04	3.0000E-06			01088	4	2	495

0.0	2.80E+06	0	0	4	01088	4	2	497
3.113E-01	1.165E-01	-4.240E-04	4.000F-06	4	1088	4	2	498
0.0	2.90E+06	0	0	4	01088	4	2	499
3.137E-01	1.166E-01	-4.440E-04	5.000E-06	4	1088	4	2	500
0.0	3.00E+06	0	0	4	01088	4	2	501
3.157E-01	1.173E-01	-4.680E-04	5.000F-06	4	1088	4	2	502
0.0	3.20E+06	0	0	4	01088	4	2	503
3.216E-01	1.186E-01	-4.810E-04	7.000F-06	4	1088	4	2	504
0.0	3.40E+06	0	0	4	01088	4	2	505
3.271E-01	1.210E-01	-5.350E-04	9.000E-06	4	1088	4	2	506
0.0	3.60E+06	0	0	4	01088	4	2	507
3.321E-01	1.252E-01	-5.480E-04	1.200F-05	4	1088	4	2	508
0.0	3.80E+06	0	0	4	01088	4	2	509
3.367E-01	1.312E-01	-5.400E-04	1.400F-05	4	1088	4	2	510
0.0	4.00E+06	0	0	4	01088	4	2	511
3.418E-01	1.381E-01	-5.780E-04	1.700F-05	4	1088	4	2	512
0.0	4.20E+06	0	0	4	01088	4	2	513
3.476E-01	1.438E-01	-5.920E-04	2.200F-05	4	1088	4	2	514
0.0	4.40E+06	0	0	4	01088	4	2	515
3.537E-01	1.495E-01	-6.100E-04	2.600F-05	4	1088	4	2	516
0.0	4.60E+06	0	0	4	01088	4	2	517
3.596E-01	1.552E-01	-6.410E-04	3.000E-05	4	1088	4	2	518
0.0	4.80E+06	0	0	4	01088	4	2	519
3.654E-01	1.604E-01	-6.740E-04	3.600E-05	4	1088	4	2	520
0.0	5.00E+06	0	0	4	01088	4	2	521
3.713E-01	1.658E-01	-6.910E-04	4.100F-05	4	1088	4	2	522
0.0	5.20E+06	0	0	4	01088	4	2	523
3.770E-01	1.707E-01	-7.030E-04	4.700E-05	4	1088	4	2	524
0.0	5.40E+06	0	0	4	01088	4	2	525
3.821E-01	1.753E-01	-7.400E-04	5.500F-05	4	1088	4	2	526
0.0	5.60E+06	0	0	4	01088	4	2	527
3.870E-01	1.800E-01	-7.320E-04	6.300E-05	4	1088	4	2	528
0.0	5.80E+06	0	0	4	01088	4	2	529
3.916E-01	1.844E-01	-7.660E-04	7.200E-05	4	1088	4	2	530
0.0	6.00E+06	0	0	4	01088	4	2	531
3.966E-01	1.884E-01	-8.190E-04	8.400E-05	4	1088	4	2	532
0.0	6.20E+06	0	0	4	01088	4	2	533
4.009E-01	1.922E-01	-8.170E-04	9.200E-05	4	1088	4	2	534
0.0	6.40E+06	0	0	4	01088	4	2	535
4.049E-01	1.953E-01	-8.200E-04	1.050F-04	4	1088	4	2	536
0.0	6.60E+06	0	0	4	01088	4	2	537
4.092E-01	1.984E-01	-8.480E-04	1.180F-04	4	1088	4	2	538
0.0	6.80E+06	0	0	4	01088	4	2	539
4.130E-01	2.011E-01	-8.700E-04	1.330E-04	4	1088	4	2	540
0.0	7.00E+06	0	0	4	01088	4	2	541
4.168E-01	2.039E-01	-8.630E-04	1.470F-04	4	1088	4	2	542
0.0	7.20E+06	0	0	4	1088	4	2	543
4.198E-01	2.062E-01	-8.140E-04	1.610E-04	4	1088	4	2	544
0.0	7.40E+06	0	0	4	1088	4	2	545
4.231E-01	2.083E-01	-7.800E-04	1.790E-04	4	1088	4	2	546
0.0	7.60E+06	0	0	4	1088	4	2	547
4.260E-01	2.106E-01	-7.500E-04	1.950E-04	4	1088	4	2	548
0.0	7.80E+06	0	0	4	1088	4	2	549
4.291E-01	2.124E-01	-6.990E-04	2.140E-04	4	1088	4	2	550
0.0	8.00E+06	0	0	4	01088	4	2	551
4.316E-01	2.140E-01	-6.410E-04	2.390E-04	4	1088	4	2	552
0.0	8.20E+06	0	0	4	01088	4	2	553
4.342E-01	2.159E-01	-5.640E-04	2.650E-04	4	1088	4	2	554
0.0	8.40E+06	0	0	4	01088	4	2	555
4.362E-01	2.175E-01	-4.670E-04	2.900E-04	4	1088	4	2	556
0.0	8.60E+06	0	0	4	01088	4	2	557
4.384E-01	2.189E-01	-3.500E-04	3.180E-04	4	1088	4	2	558

0.0	8.80E+06	0	0	4	01088 4	2	559
4.403E-01	2.202E-01	-2.490E-04	3.480E-04	4	1088 4	2	560
0.0	9.00E+06	0	0	4	01088 4	2	561
4.424E-01	2.216E-01	-1.320E-04	3.800E-04	4	1088 4	2	562
0.0	9.20E+06	0	0	4	01088 4	2	563
4.445E-01	2.223E-01	-1.200E-05	4.110E-04	4	1088 4	2	564
0.0	9.40E+06	0	0	4	01088 4	2	565
4.463E-01	2.234E-01	1.110E-04	4.480E-04	4	1088 4	2	566
0.0	9.60E+06	0	0	4	01088 4	2	567
4.482E-01	2.245E-01	2.700E-04	4.840E-04	4	1088 4	2	568
0.0	9.80E+06	0	0	4	01088 4	2	569
4.498E-01	2.254E-01	4.340E-04	5.220E-04	4	1088 4	2	570
0.0	1.00E+07	0	0	4	01088 4	2	571
4.517E-01	2.265E-01	6.160E-04	5.600E-04	4	1088 4	2	572
0.0	1.05E+07	0	0	4	01088 4	2	573
4.553E-01	2.287E-01	1.128E-03	6.820E-04	4	1088 4	2	574
0.0	1.10E+07	0	0	4	01088 4	2	575
4.589E-01	2.310E-01	1.609E-03	7.830E-04	4	1088 4	2	576
0.0	1.15E+07	0	0	4	01088 4	2	577
4.621E-01	2.326E-01	2.313E-03	9.250E-04	4	1088 4	2	578
0.0	1.20E+07	0	0	4	01088 4	2	579
4.651E-01	2.340E-01	3.169E-03	1.083E-03	4	1088 4	2	580
0.0	1.25E+07	0	0	4	01088 4	2	581
4.674E-01	2.351E-01	4.082E-03	1.231E-03	4	1088 4	2	582
0.0	1.30E+07	0	0	4	01088 4	2	583
4.701E-01	2.362E-01	5.001E-03	1.429E-03	4	1088 4	2	584
0.0	1.35E+07	0	0	4	01088 4	2	585
4.725E-01	2.374E-01	6.139E-03	1.600E-03	4	1088 4	2	586
0.0	1.40E+07	0	0	4	01088 4	2	587
4.750E-01	2.388E-01	7.284E-03	1.781E-03	4	1088 4	2	588
0.0	1.45E+07	0	0	4	01088 4	2	589
4.775E-01	2.402E-01	8.515E-03	1.971E-03	4	1088 4	2	590
0.0	1.50E+07	0	0	4	01088 4	2	591
4.801E-01	2.418E-01	9.809E-03	2.171E-03	4	1088 4	2	592
0.0	0.0	0	0	0	01088 4	0	593
0.0	0.0	0	0	0	01088 0	0	594
2000.0	3.96822	0	0	0	01088 7	4	595
0.0	0.0	0	0	12	11088 7	4	596
0.0	59.3	3.96822	1.5	0.0	0.01088 7	4	597
1.0	0.7238	3.96822	0.0	0.0	0.01088 7	4	598
0.0	0.0	0	0	0	01088 7	0	599
0.0	0.0	0	0	0	01088 0	0	600
0.0	0.0	0	0	0	0	0	601

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